

# Saddle point problems in liquid crystal modelling

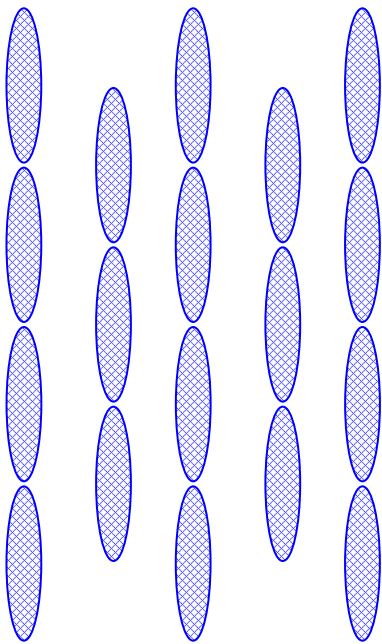
Alison Ramage  
Dept of Mathematics  
University of Strathclyde  
Glasgow, Scotland



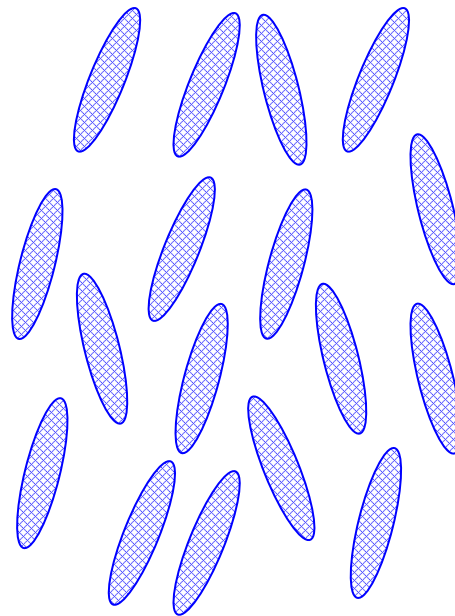
Eugene C. Gartland, Jr.  
Dept of Mathematics  
Kent State University  
Ohio, USA

# Liquid Crystals

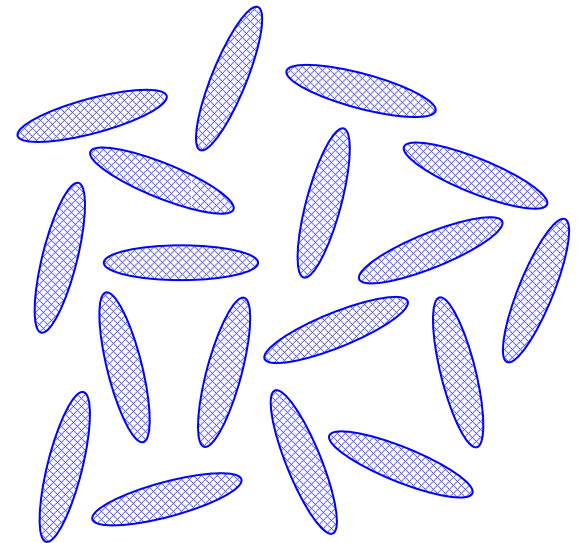
- occur between solid crystal and isotropic liquid states



solid



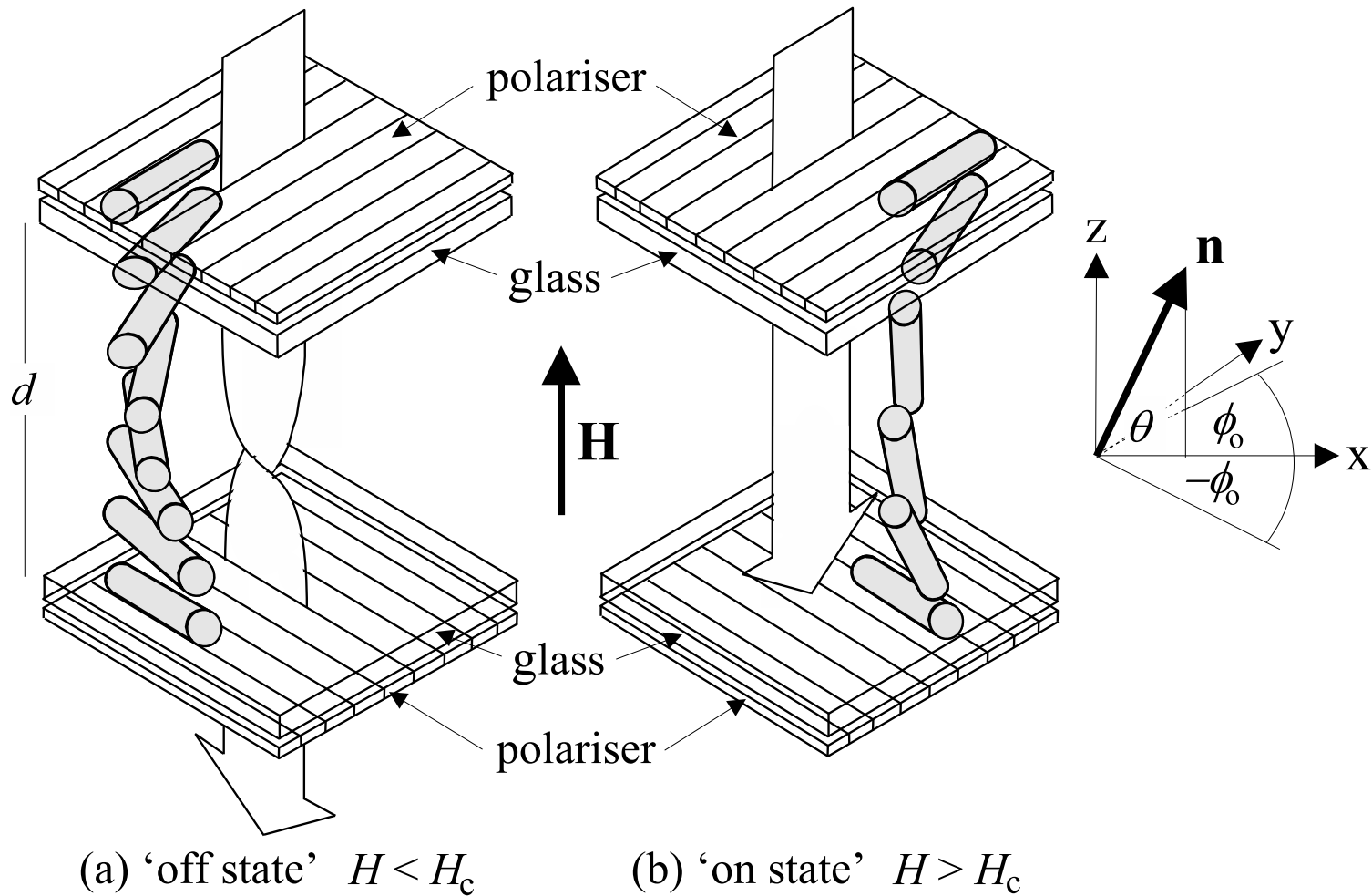
liquid crystal



liquid

- may have different **equilibrium** configurations
- **switch** between stable states by altering applied voltage, magnetic field, boundary conditions, . . .

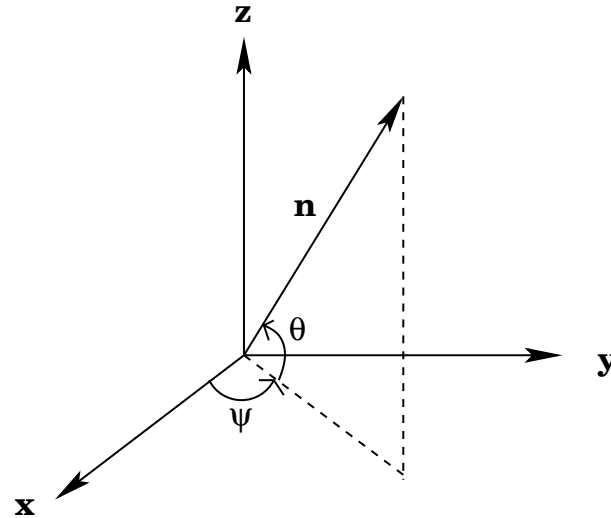
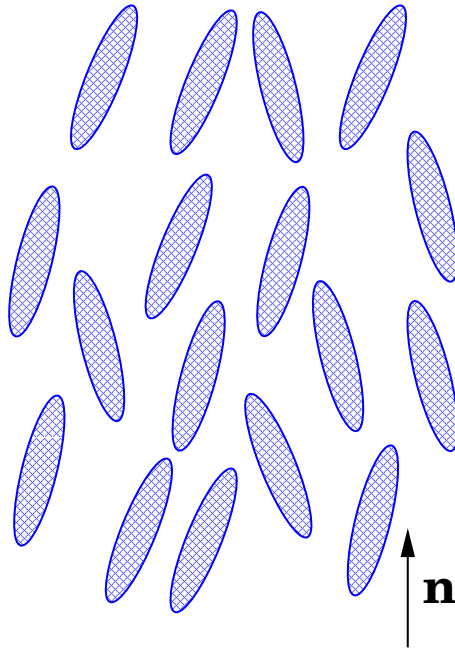
# Liquid Crystal Displays



twisted nematic device

*Static and Dynamic Continuum Theory of Liquid Crystals,*  
*Iain W. Stewart (2004)*

# Modelling: Director-based models



- **director**: average direction of molecular alignment

unit vector

$$\mathbf{n} = (\cos \theta \cos \psi, \cos \theta \sin \psi, \sin \theta)$$

- **order parameter**: measure of orientational order

$$S = \frac{1}{2} \langle 3 \cos^2 \theta_m - 1 \rangle$$

# Finding Equilibrium Configurations

- minimise the **free energy density**

$$\mathcal{F} = \int_V F_{bulk}(\theta, \psi, \nabla\theta, \nabla\psi) + \int_{\mathcal{S}} F_{surface}(\theta, \phi) d\mathcal{S}$$

$$F_{bulk} = F_{elastic} + F_{electrostatic}$$

- if fixed boundary conditions are applied, surface energy term can be ignored
- solutions with **least** energy are physically relevant
- use calculus of variations: **Euler-Lagrange equations**

# Elastic Energy

- Frank-Oseen elastic energy

$$\begin{aligned} F_{elastic} &= \frac{1}{2}K_1(\nabla \cdot \mathbf{n})^2 + \frac{1}{2}K_2(\mathbf{n} \cdot \nabla \times \mathbf{n})^2 \\ &+ \frac{1}{2}K_3(\mathbf{n} \times \nabla \times \mathbf{n})^2 \\ &+ \frac{1}{2}(K_2 + K_4)\nabla \cdot [(\mathbf{n} \cdot \nabla)\mathbf{n} - (\nabla \cdot \mathbf{n})\mathbf{n}] \end{aligned}$$

- Frank elastic constants

$K_1$	splay
$K_2$	twist
$K_3$	bend
$K_2 + K_4$	saddle-splay

# One-Constant Approximation

- set

$$K = K_1 = K_2 = K_3, \quad K_4 = 0$$

- vector identities

$$(\nabla \times \mathbf{n})^2 = (\mathbf{n} \cdot \nabla \times \mathbf{n})^2 + (\mathbf{n} \times \nabla \times \mathbf{n})^2$$

$$\nabla(\mathbf{n} \cdot \mathbf{n}) = 0$$

$$[(\nabla \cdot \mathbf{n})^2 + (\nabla \times \mathbf{n})^2] + \nabla \cdot [(\mathbf{n} \cdot \nabla)\mathbf{n} - (\nabla \cdot \mathbf{n})\mathbf{n}] = \|\nabla \mathbf{n}\|^2$$

- elastic energy  $F_{elastic} = \frac{1}{2}K \|\nabla \mathbf{n}\|^2$

# Electrostatic Energy

- applied electric field  $\mathbf{E}$  of magnitude  $E$
- electrostatic energy

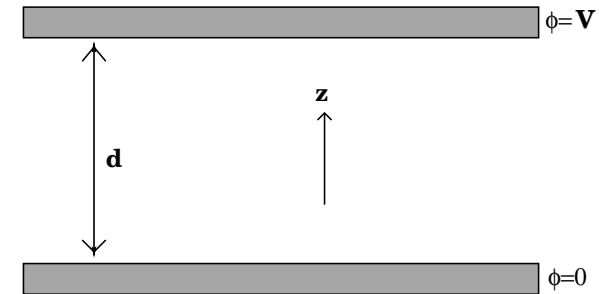
$$F_{electrostatic} = -\frac{1}{2}\epsilon_0\epsilon_{\perp}E^2 - \frac{1}{2}\epsilon_0\epsilon_a(\mathbf{n} \cdot \mathbf{E})^2$$

- dielectric anisotropy  $\epsilon_a = \epsilon_{\parallel} - \epsilon_{\perp}$
- permittivity of free space  $\epsilon_0$



# Model Problem: Twisted Nematic Device

- two parallel plates distance  $d$  apart



- **strong anchoring** parallel to plate surfaces ( $\mathbf{n}$  fixed)
- rotate one plate through  $\pi/2$  radians
- electric field  $\mathbf{E} = (0, 0, E(z))$ , voltage  $V$

# Equilibrium Equations 1

- equilibrium equations on  $z \in [0, d]$

$$F = \frac{1}{2} \int_0^d \left\{ K \|\nabla \mathbf{n}\|^2 - \epsilon_0 \epsilon_{\perp} E^2 - \epsilon_0 \epsilon_a (\mathbf{n} \cdot \mathbf{E})^2 \right\} dz$$

- director  $\mathbf{n} = (u, v, w), \quad |\mathbf{n}| = 1$
- constraint applied via Lagrange multipliers  $\lambda$
- electric potential  $\phi: \quad E = \frac{d\phi}{dz}$
- unknowns  $u, v, w, \phi, \lambda$

# Equilibrium Equations 2

- nondimensionalised equilibrium equations on  $z \in [0, 1]$

$$F = \frac{1}{2} \int_0^1 \left[ (u_z^2 + v_z^2 + w_z^2) - \alpha^2 \pi^2 (\beta + w^2) \phi_z^2 - \lambda(u^2 + v^2 + w^2 - 1) \right] dz$$

- dimensionless parameters

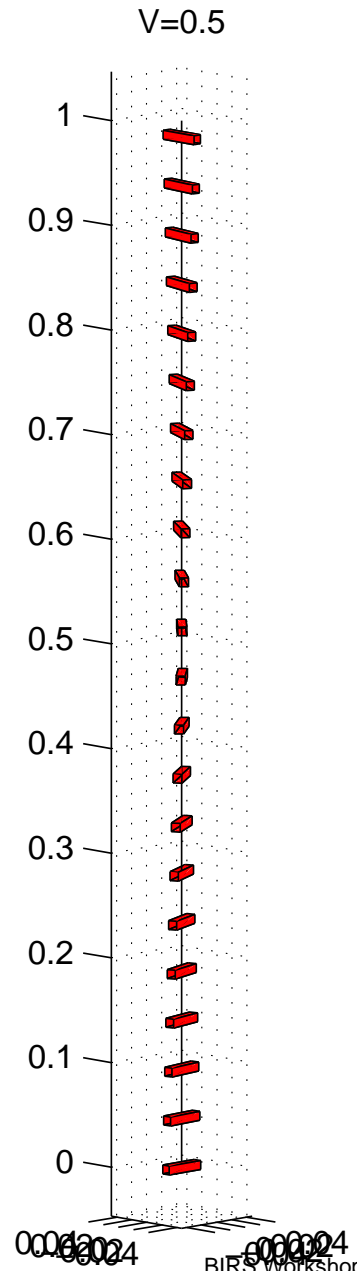
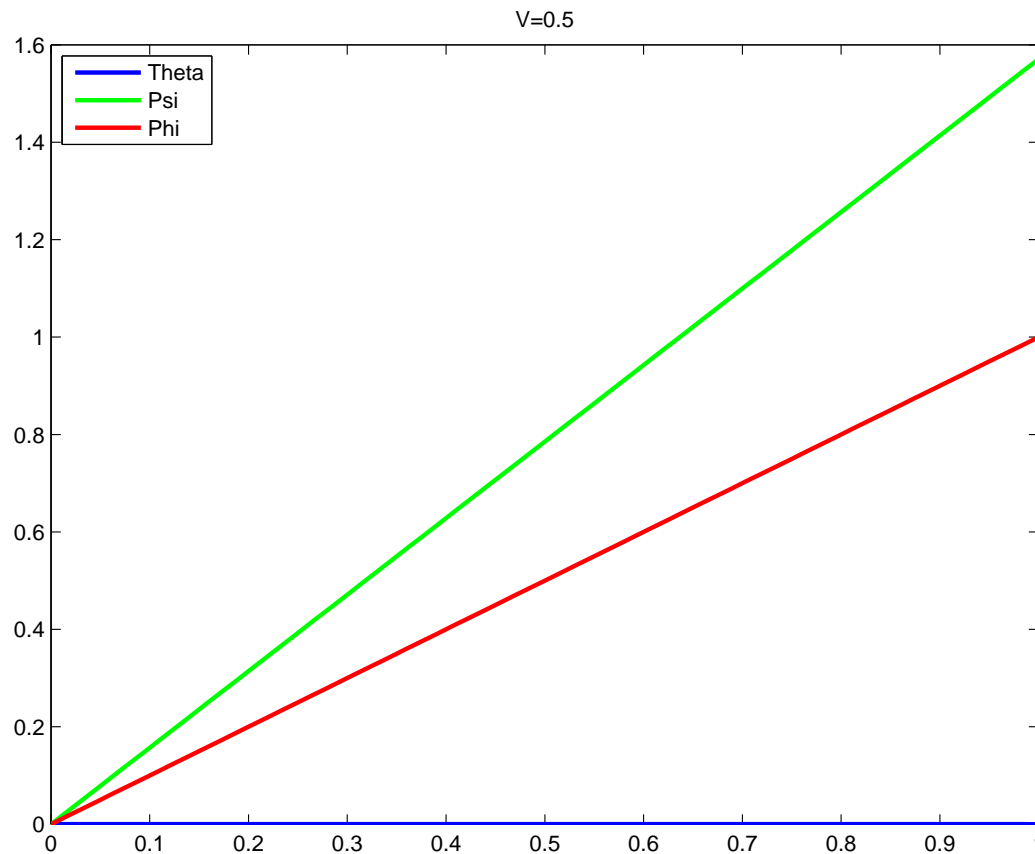
$$\alpha^2 = \frac{\epsilon_0 \epsilon_a V^2}{K \pi^2}, \quad \beta = \frac{\epsilon_{\perp}}{\epsilon_a}$$

- boundary conditions:

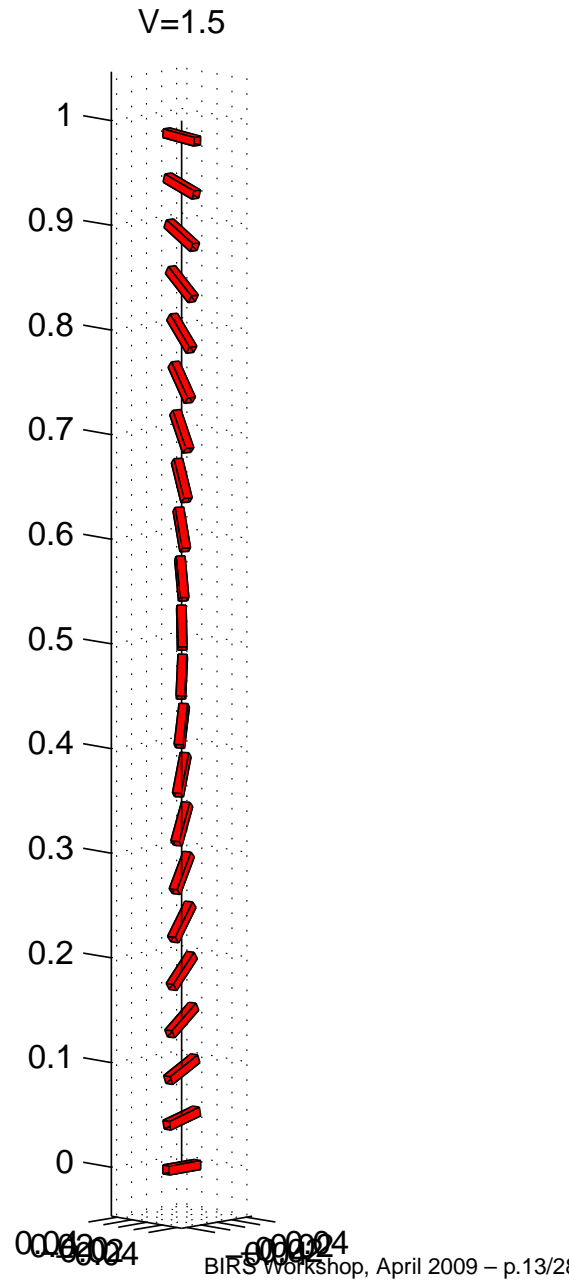
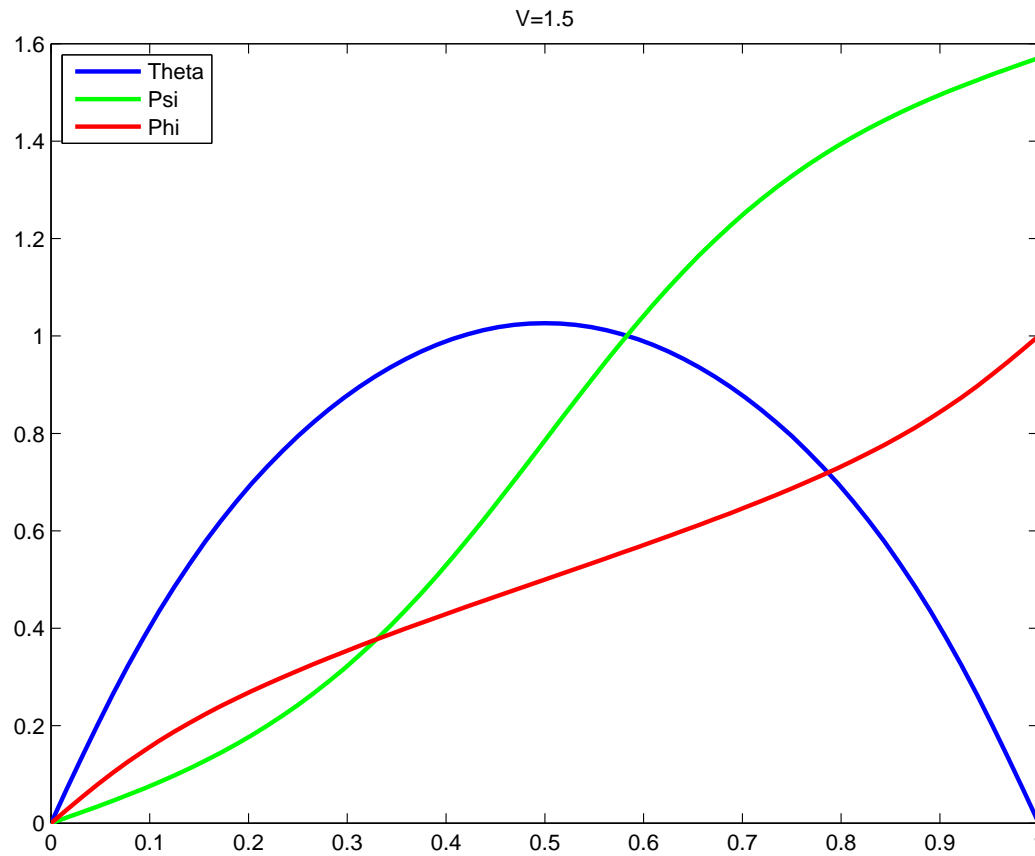
$$\text{at } z = 0: \theta = \psi = \phi = 0, \quad \text{at } z = 1: \theta = 0, \psi = \frac{\pi}{2}, \phi = 1$$

# Off State

$$\theta(z) \equiv 0, \quad \psi(z) = \frac{\pi z}{2}, \quad \phi(z) = z$$



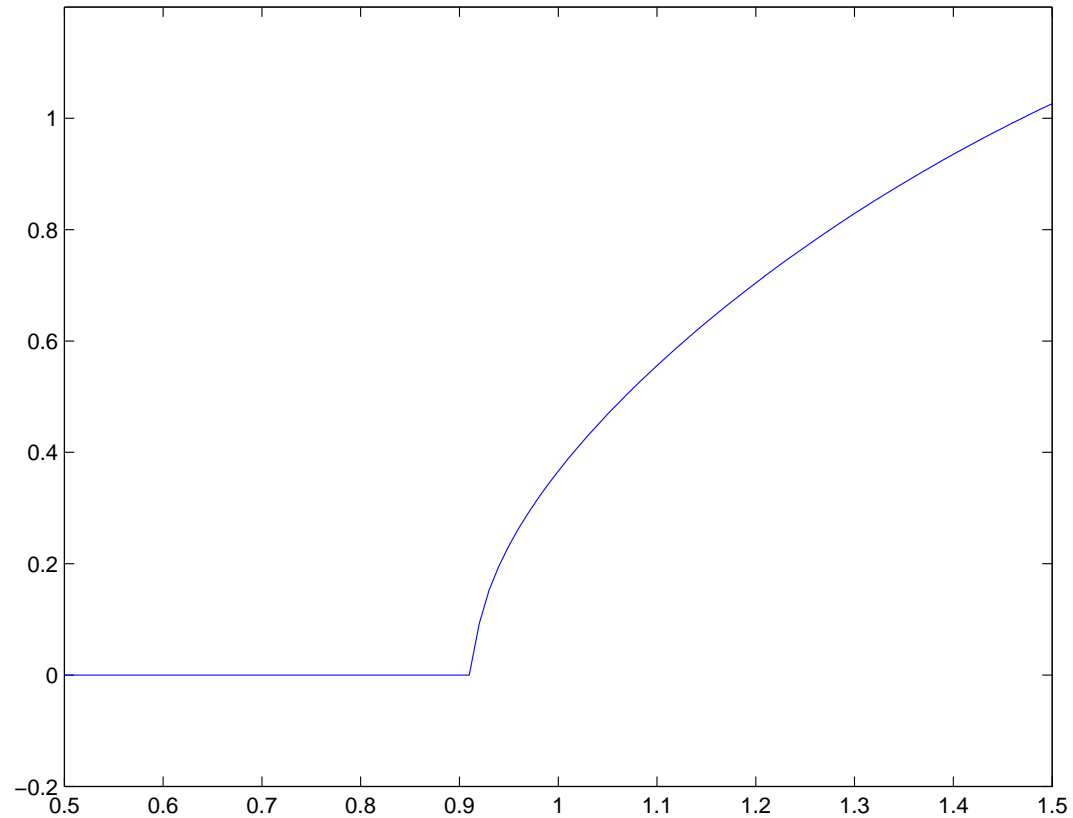
# On State



# Critical Voltage

- switching occurs at

$$V_c = \frac{\pi}{2} \sqrt{\frac{3K}{\epsilon_0 \epsilon_a}}$$



# Discrete Free Energy

- grid of  $N + 1$  points  $z_k$  a distance  $\Delta z$  apart
- approximate integral by **mid-point** rule

$$\begin{aligned} F &\simeq \frac{\Delta z}{2} \sum_{k=0}^{N-1} \left\{ \left[ \frac{u_{k+1} - u_k}{\Delta z} \right]^2 + \left[ \frac{v_{k+1} - v_k}{\Delta z} \right]^2 + \left[ \frac{w_{k+1} - w_k}{\Delta z} \right]^2 \right. \\ &- \alpha^2 \pi^2 \left( \beta + \left[ \frac{w_k^2 + w_{k+1}^2}{2} \right] \right) \left[ \frac{\phi_{k+1} - \phi_k}{\Delta z} \right]^2 \\ &\left. - \lambda_k \left[ \frac{u_k^2 + u_{k+1}^2}{2} + \frac{v_k^2 + v_{k+1}^2}{2} + \frac{w_k^2 + w_{k+1}^2}{2} - 1 \right] \right\} \end{aligned}$$

# Euler-Lagrange Equations

- set  $\frac{\partial F}{\partial u_k}, \frac{\partial F}{\partial v_k}, \frac{\partial F}{\partial w_k}, \frac{\partial F}{\partial \phi_k}, \frac{\partial F}{\partial \lambda_k}$  equal to zero

- solve  $\nabla \mathbf{F}(\mathbf{x}) = \mathbf{0}$  for  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \phi, \lambda]$   
 $N + 1$  gridpoints  $\Rightarrow n = N - 1$  unknowns

- use Newton's method: solve

$$\nabla^2 \mathbf{F}(\mathbf{x}_j) \cdot \delta \mathbf{x}_j = -\nabla \mathbf{F}(\mathbf{x}_j)$$

- linear system, coefficient matrix is **Hessian**  $\nabla^2 \mathbf{F}(\mathbf{x}_j)$

$$\nabla^2 \mathbf{F} = \begin{bmatrix} \nabla_{\mathbf{nn}}^2 \mathbf{F} & \nabla_{\mathbf{n}\phi}^2 \mathbf{F} & \nabla_{\mathbf{n}\lambda}^2 \mathbf{F} \\ \nabla_{\phi\mathbf{n}}^2 \mathbf{F} & \nabla_{\phi\phi}^2 \mathbf{F} & \nabla_{\phi\lambda}^2 \mathbf{F} \\ \nabla_{\lambda\mathbf{n}}^2 \mathbf{F} & \nabla_{\lambda\phi}^2 \mathbf{F} & \nabla_{\lambda\lambda}^2 \mathbf{F} \end{bmatrix}$$



# Hessian Components 1

- matrix notation:  $\nabla_{\mathbf{nn}}^2 \mathbf{F} = A$

$$A = \begin{bmatrix} \nabla_{\mathbf{uu}}^2 \mathbf{F} & 0 & 0 \\ 0 & \nabla_{\mathbf{vv}}^2 \mathbf{F} & 0 \\ 0 & 0 & \nabla_{\mathbf{ww}}^2 \mathbf{F} \end{bmatrix} = \begin{bmatrix} A_{uu} & 0 & 0 \\ 0 & A_{vv} & 0 \\ 0 & 0 & A_{ww} \end{bmatrix}$$

- $A_{uu}$ ,  $A_{vv}$  and  $A_{ww}$  are  $n \times n$  **symmetric tridiagonal** blocks

- $A_{uu} = A_{vv} = \frac{1}{\Delta z} \text{tridiag}(-1, 2 - \Delta z^2 \lambda_j, -1)$

- $A_{ww} = \frac{1}{\Delta z} \text{tridiag}(-1, 2 - \Delta z^2 \lambda_j - \gamma_j, -1)$

$$\gamma_j = \frac{\alpha^2 \pi^2}{2} [(\phi_{j+1} - \phi_j)^2 + (\phi_j - \phi_{j-1})^2]$$

# Eigenvalues of $A$

- at first Newton step (initial linear  $\phi$ ,  $\lambda_j = 1$ ) block matrices are **Toeplitz**
- find eigenvalues using Fourier analysis
- $\sigma_{\min}(A_{uu}) = \sigma_{\min}(A_{vv}) \simeq \Delta z(\pi^2 - \lambda_1) > 0$   
 $A_{uu}$  and  $A_{vv}$  are initially positive definite
- $\sigma_{\min}(A_{ww}) \simeq \Delta z(\pi^2(1 - \alpha^2) - \lambda_1)$   
 $A_{ww}$  is initially positive definite iff  $V < \frac{2}{\sqrt{3}}V_c$
- at subsequent Newton iterations,  $A_{uu}$ ,  $A_{vv}$ ,  $A_{ww}$  may all be **indefinite**
- number of negative eigenvalues increases with  $V$

# Hessian Components 2

- matrix notation:  $\nabla_{\mathbf{n}\lambda}^2 \mathbf{F} = B$

- the  $3n \times n$  matrix  $B$  has structure

$$B = \Delta z \begin{bmatrix} B_u \\ B_v \\ B_w \end{bmatrix}, \quad \begin{aligned} B_u &= \text{diag}(\mathbf{u}) \\ B_v &= \text{diag}(\mathbf{v}) \\ B_w &= \text{diag}(\mathbf{w}) \end{aligned}$$

- $\text{rank}(B^T) = n$

- $B^T B = \Delta z^2 I_n$

- information available about basis for nullspace of  $B^T$

# Hessian Components 3

- matrix notation:  $\nabla_{\phi\phi}^2 \mathbf{F} = -C$
- the  $n \times n$  matrix  $C$  is **symmetric** and **tridiagonal**
- $C = \frac{1}{\Delta z} \text{tridiag}(-a_{j-\frac{1}{2}}, a_{j-\frac{1}{2}} + a_{j+\frac{1}{2}}, -a_{j+\frac{1}{2}})$

$$a_{j-\frac{1}{2}} = \alpha^2 \pi^2 \left( \beta + \frac{1}{2} (w_{j-1}^2 + w_j^2) \right) > 0$$

$$a_{j+\frac{1}{2}} = \alpha^2 \pi^2 \left( \beta + \frac{1}{2} (w_j^2 + w_{j+1}^2) \right) > 0$$

- diagonally dominant with positive real diagonal entries

$C$  is positive definite

# Hessian Components 4

- matrix notation:  $\nabla_{\mathbf{n}\phi}^2 \mathbf{F} = D$

$$D = \Delta z \begin{bmatrix} 0 \\ 0 \\ \mu D_w \end{bmatrix}, \quad \mu = \frac{\alpha^2 \pi^2}{\Delta z}$$

- the  $n \times n$  matrix  $D_w$  is **tridiagonal**

$$D_w = \text{diag}(\mathbf{w}) \text{tridiag}(\phi_j - \phi_{j-1}, \phi_{j-1} - 2\phi_j + \phi_{j+1}, \phi_j - \phi_{j+1})$$

- $D_w$  has **complex** eigenvalues (including one zero)
- $\text{rank}(D) = n - 1$

# Four Saddle-Point Problems

- for unknown vector ordered as  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \phi, \lambda]$

$$H = \left[ \begin{array}{c|cc} A & D & B \\ \hline D^T & -C & 0 \\ B^T & 0 & 0 \end{array} \right]$$

$$H = \left[ \begin{array}{cc|c} A & D & B \\ \hline D^T & -C & 0 \\ B^T & 0 & 0 \end{array} \right]$$

- for unknown vector ordered as  $\mathbf{x} = [\mathbf{u}, \mathbf{v}, \mathbf{w}, \lambda, \phi]$

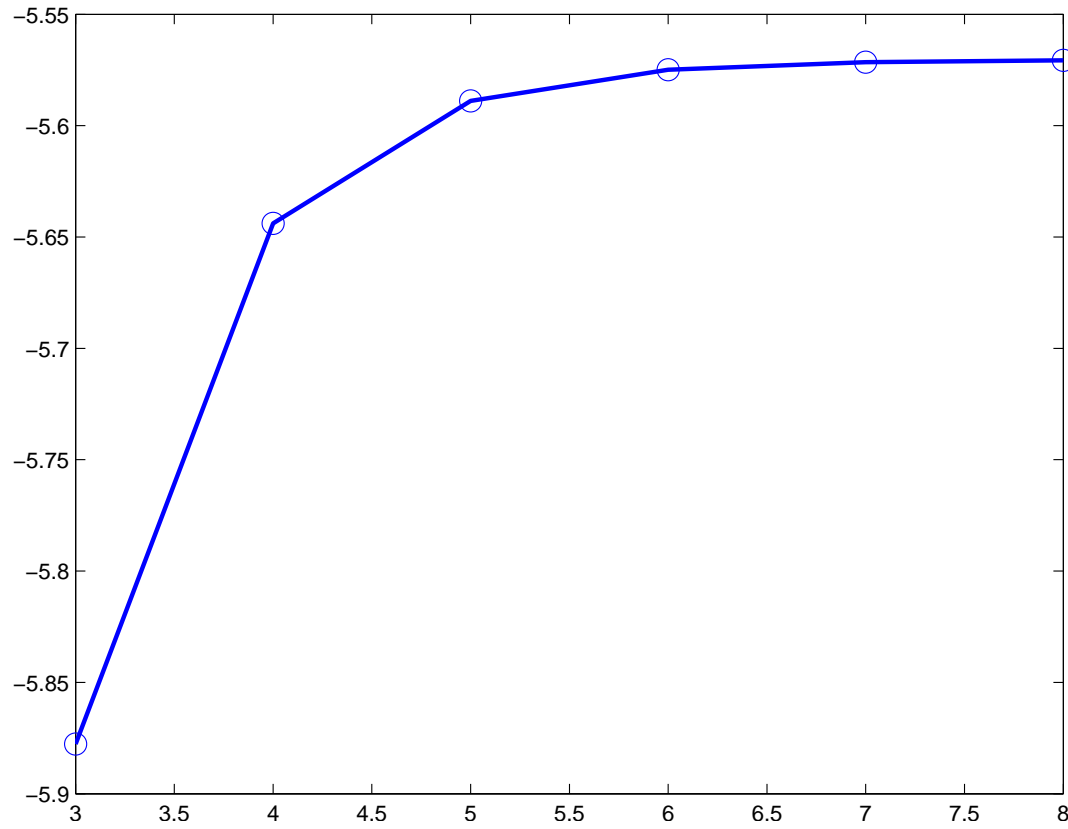
$$H = \left[ \begin{array}{c|cc} A & B & D \\ \hline B^T & 0 & 0 \\ D^T & 0 & -C \end{array} \right]$$

$$H = \left[ \begin{array}{cc|c} A & B & D \\ \hline B^T & 0 & 0 \\ D^T & 0 & -C \end{array} \right]$$

**double** saddle-point structure

# Iterative Solution

- outer iteration: **Newton's method**  $\text{tol}=1e-4$
- inner iteration: **MINRES**  $\text{tol}=1e-4$
- check accuracy by calculating energy of final solution



# Matrix Conditioning

- eigenvalues of  $H$  lie in  $[\lambda_{\min}, \lambda_s] \cup [\lambda_{s+1}, \lambda_{\max}]$
- estimate of matrix conditioning:

$N$	condest	$\lambda_{\min}(H)$	$\lambda_s(H)$	$\lambda_{s+1}(H)$	$\lambda_{\max}(H)$
8	1.64e+6	-6.68e+2	-5.40e-4	1.88e-1	3.07e+1
16	2.58e+7	-1.44e+3	-6.26e-5	2.19e-1	6.33e+1
32	4.09e+8	-2.98e+3	-7.68e-6	1.28e-1	1.28e+2
64	6.51e+9	-6.07e+3	-9.56e-7	6.60e-2	2.56e+2
128	1.04e+11	-1.23e+4	-1.20e-7	3.33e-2	5.12e+2
256	1.66e+12	-2.46e+4	-1.50e-8	1.67e-2	1.03e+3
	$O(N^4)$	$O(N)$	$O(N^{-3})$	$O(N^{-1})$	$O(N)$



# Diagonal Preconditioning

$$H = \begin{bmatrix} A & D & B \\ D^T & -C & 0 \\ B^T & 0 & 0 \end{bmatrix}$$

$$\mathcal{D} = \begin{bmatrix} D_A & 0 & 0 \\ 0 & D_C & 0 \\ 0 & 0 & \Delta z I \end{bmatrix} \quad \begin{aligned} D_A &= |\text{diag}(A)| \\ D_C &= \text{diag}(C) \end{aligned}$$

- estimated condition of  $P^{-1}H$  is  $O(N^2)$

$$\lambda_{\min} = -2, \quad \lambda_s = O(N^{-2}), \quad \lambda_{s+1} = O(N^{-2}), \quad \lambda_{\max} = 2.$$

# Constraint-type Preconditioning

$$H = \left[ \begin{array}{cc|c} A & B & D \\ B^T & 0 & 0 \\ \hline D^T & 0 & -C \end{array} \right]$$

$$C_1 = \left[ \begin{array}{cc|c} D_A & 0 & D \\ 0 & \Delta z I & 0 \\ \hline D^T & 0 & -C \end{array} \right], \quad C_2 = \left[ \begin{array}{cc|c} A & 0 & D \\ 0 & \Delta z I & 0 \\ \hline D^T & 0 & -C \end{array} \right]$$

- Projected Preconditioned Conjugate Gradients  
Dollar et al. (2006)

# Iteration Counts

- iteration counts at **first** Newton step

$N$	8	16	32	64	128	256
$\mathcal{D}$	15	40	117	382	1293	5126
$C_1$	13	25	50	98	195	387
$C_2$	7	9	8	9	7	8

- iteration counts at **last** Newton step

$N$	8	16	32	64	128	256
$\mathcal{D}$	37	134	414	1617	7466	34755
$C_1$	22	55	226	635	2259	7166
$C_2$	6	14	23	43	65	114

# Other methods?

- block tridiagonal?
- more sophisticated constraint preconditioning?
- Schur complement approximation?
- augmented Lagrangian methods?
- inner/outer iteration?
- connection with harmonic maps?