

Multigrid Solution of Discrete Convection-Diffusion Equations

Alison Ramage
Dept of Mathematics
University of Strathclyde
Glasgow, Scotland



Howard Elman
Dept of Computer Science
University of Maryland
College Park, MD, USA



Overview

- background
 - convection-diffusion problems
 - multigrid methods
- practical multigrid issues
 - approximation and smoothing properties
 - convergence analysis
- model problem Fourier analysis
 - matrix transformation
 - comparison with semiperiodic problem
 - implications for Dirichlet problems

Convection-Diffusion in 2D

$$\begin{aligned}-\epsilon \nabla^2 u(x, y) + \mathbf{w} \cdot \nabla u(x, y) &= f(x, y) \quad \text{in} \quad \Omega \subset \mathbb{R}^2 \\ u(x, y) &= g \quad \text{on} \quad \partial\Omega\end{aligned}$$

divergence-free convective velocity ('wind') \mathbf{w}

diffusion parameter $\epsilon \ll 1$

discretisation parameter h

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diffusion parameter $\epsilon \ll 1$

discretisation parameter h

mesh Péclet number $P_h = \frac{\|\mathbf{w}\| h}{2\epsilon}$

Boundary Layers and Oscillations

- Galerkin finite element method

$$\epsilon(\nabla u_h, \nabla v_h) + (\mathbf{w} \cdot \nabla u_h, v_h) = (f, v_h) \quad \forall v_h \in V_h$$

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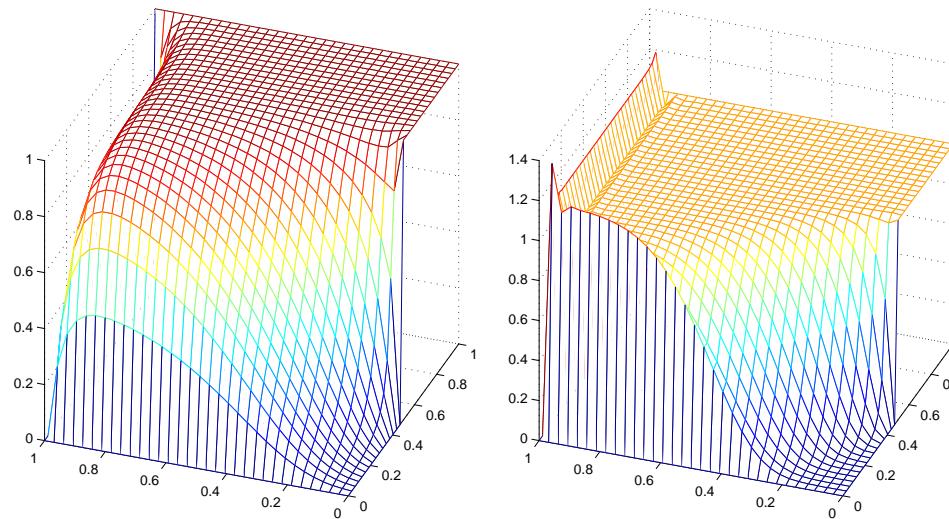
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- solution features:
exponential and **characteristic** boundary layers
- **oscillations** observed in discrete solutions for $P_h > 1$



$$P_e = 0.5$$

$$P_e = 2$$

Streamline Diffusion Method

streamline diffusion FEM, square bilinear elements

$$\begin{aligned}\epsilon(\nabla u_h, \nabla v_h) + (\mathbf{w} \cdot \nabla u_h, v_h) + \frac{\delta h}{\|\mathbf{w}\|} (\mathbf{w} \cdot \nabla u_h, \mathbf{w} \cdot \nabla v_h) \\ = (f, v_h) + \frac{\delta h}{\|\mathbf{w}\|} (f, \mathbf{w} \cdot \nabla v_h) \quad \forall v_h \in V_h\end{aligned}$$

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- $P_h \leq 1$: $\delta = 0$ Galerkin FEM

- $P_h > 1$: $\delta = \frac{1}{2} - \frac{\epsilon}{h}$ Streamline Diffusion

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 - rapidly reduces error components in subspace 2
- recursive process on nested grids
- optimal in the sense of obtaining convergence rate independent of h

Issues for Convection-Diffusion

- approximation: choice of discretisation
 - oscillations on coarser grids?
 - grid transfer operators?

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- multigrid can be implemented effectively for convection-diffusion problems

Convergence Analysis

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- ideas for convection-diffusion less well-developed

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- standard Poisson-type convergence analysis fails
- ideas for convection-diffusion less well-developed
- various successful approaches
 - perturbation arguments
Bank (1981), Bramble, Pasciak and Xu (1988),
Mandel (1986), Wang (1993)
 - matrix-based methods
Reusken (2002), Olyshanskii and Reusken (2002)

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direct discretisation on coarse grid

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- restriction: transpose of prolongation P^T

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- prolongation: bilinear interpolation P
- restriction: transpose of prolongation P^T
- smoothing: line Gauss-Seidel S_A
- ν steps of pre-smoothing, no post-smoothing

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- convergence?

$$\|\mathbf{e}_k\| \leq \|M\|^k \|\mathbf{e}_0\|$$

convergence if $\|M\| < 1$

Two-Grid Convergence Analysis

AIM: find an upper bound for

$$\|M\|_2 = \|(I - PA_c^{-1}P^T A_f)S_A^\nu\|_2$$

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- Approach 1: write

$$M = (\textcolor{red}{A_f^{-1}} - PA_c^{-1}P^T)(\textcolor{green}{A_f S_A^\nu}) = M_A M_S$$

and bound $\|M_A\|_2$, $\|M_S\|_2$ separately

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- Approach 2: bound $\|M\|_2$ directly

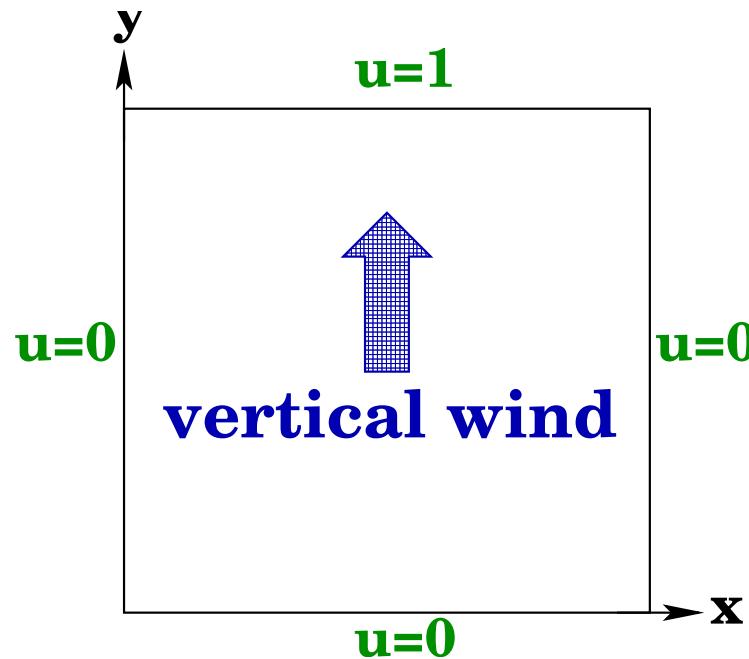
Model Problem

grid-aligned flow with vertical wind and $f = 0$

$$-\epsilon \nabla^2 u(x, y) + (0, 1) \cdot \nabla u(x, y) = 0$$

Dirichlet boundary conditions

square bilinear elements



Computational Molecule

parameters h, ϵ, δ

$$M_2 : -\frac{1}{12}[(2\delta-1)h+4\epsilon] \quad -\frac{1}{3}[(2\delta-1)h+\epsilon] \quad -\frac{1}{12}[(2\delta-1)h+4\epsilon]$$



$$M_1 : \quad \frac{1}{3}(\delta h - \epsilon) \quad \leftarrow \quad \frac{4}{3}(\delta h + 2\epsilon) \quad \rightarrow \quad \frac{1}{3}(\delta h - \epsilon)$$



$$M_3 : -\frac{1}{12}[(2\delta+1)h+4\epsilon] \quad -\frac{1}{3}[(2\delta+1)h+\epsilon] \quad -\frac{1}{12}[(2\delta+1)h+4\epsilon]$$

symmetric stencil

Coefficient Matrix

$$A = \begin{bmatrix} M_1 & M_2 & & & 0 \\ M_3 & M_1 & M_2 & & \\ & \ddots & \ddots & \ddots & \\ & & M_3 & M_1 & M_2 \\ 0 & & & M_3 & M_1 \end{bmatrix}$$

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eigenvectors and eigenvalues:

$$M_1 \mathbf{v}_j = \lambda_j \mathbf{v}_j, \quad \lambda_j = m_{1c} + 2m_{1r} \cos \frac{j\pi}{N}$$

$$M_2 \mathbf{v}_j = \sigma_j \mathbf{v}_j, \quad \sigma_j = m_{2c} + 2m_{2r} \cos \frac{j\pi}{N}$$

$$M_3 \mathbf{v}_j = \gamma_j \mathbf{v}_j, \quad \gamma_j = m_{3c} + 2m_{3r} \cos \frac{j\pi}{N}$$

$$\mathbf{v}_j = \sqrt{\frac{2}{N}} \left[\sin \frac{j\pi}{N}, \quad \sin \frac{2j\pi}{N}, \quad \dots, \sin \frac{(N-1)j\pi}{N} \right]^T$$

Transformation: Coefficient Matrix (1)

N_f^2 elements, n_f^2 unknowns ($n_f = N_f - 1$)

$$\hat{V}_f = [\mathbf{v}_1 \mathbf{v}_2 \dots \mathbf{v}_{n_f}], \quad V_f = \text{diag}(\hat{V}_f, \dots, \hat{V}_f)$$

$$M_1 \hat{V}_f = \hat{V}_f \Lambda, \quad M_2 \hat{V}_f = \hat{V}_f \Sigma, \quad M_3 \hat{V}_f = \hat{V}_f \Gamma$$

$$V_f^T A_f V_f = \hat{T}_f = \begin{bmatrix} \Lambda & \Sigma & & & 0 \\ \Gamma & \Lambda & \Sigma & & \\ & \ddots & \ddots & \ddots & \\ & & \Gamma & \Lambda & \Sigma \\ 0 & & & \Gamma & \Lambda \end{bmatrix}$$

Transformation: Coefficient Matrix (2)

permute into tridiagonal form:

$$\Pi_f^T \hat{T}_f \Pi_f = T_f = \begin{bmatrix} T_1 & & & & 0 \\ & T_2 & & & \\ & & \ddots & & \\ & & & T_{n_f-1} & \\ 0 & & & & T_{n_f} \end{bmatrix}$$

$$T_j = \text{tridiag}(\gamma_j, \lambda_j, \sigma_j)$$

$$A_f = Q_f T_f Q_f^T \quad Q_f = V_f \Pi_f$$

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coarse grid: $A_c = Q_c T_c Q_c^T \quad Q_c = V_c \Pi_c$

Transformation: Smoothing Matrix

block matrix splitting: $A_f = D_A - L_A - U_A$

Gauss-Seidel smoothing matrix:

$$S_A = (D_A - L_A)^{-1}U_A = I - (D_A - L_A)^{-1}A_f$$

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transformation:

$$S_A = Q_f S_T Q_f^T$$

where $S_T = I - (D_T - L_T)^{-1}T_f$ is block-diagonal

Transformation: Prolongation Matrix

2D prolongation matrix: $P = L \otimes L$

$$L^T = \begin{bmatrix} \frac{1}{2} & 1 & \frac{1}{2} & & & \\ & \frac{1}{2} & 1 & \frac{1}{2} & & \\ & & \ddots & \ddots & \ddots & \\ & & & \frac{1}{2} & 1 & \frac{1}{2} \end{bmatrix}$$

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transformation: $Q_f = (I_f \otimes \hat{V}_f) \Pi_f, \quad Q_c = (I_c \otimes \hat{V}_c) \Pi_c$

$$\bar{P} = Q_f^T P Q_c = \mathcal{A}^T \otimes L$$

$$\mathcal{A} = \begin{bmatrix} \alpha_1 & & & & & \alpha_{N_f-1} \\ & \alpha_2 & & & & \alpha_{N_f-2} \\ & & \ddots & & & \ddots \\ & & & \alpha_{n_c} & 0 & \alpha_{N_f-n_c} \end{bmatrix}$$

Transformation: Iteration Matrix (1)

$$\begin{aligned} M &= (I - PA_c^{-1}P^T A_f) S_A^\nu \\ &= (I - PQ_c T_c^{-1} Q_c^T P^T Q_f T_f Q_f^T) S_A^\nu \\ &= Q_f (I - \bar{P} T_c^{-1} \bar{P}^T T_f) Q_f^T (Q_f S_T Q_f^T)^\nu \\ &= Q_f \left(I - \bar{P} T_c^{-1} \bar{P}^T T_f \right) S_T^\nu Q_f^T \\ \Rightarrow M &= Q_f \bar{M} Q_f^T \end{aligned}$$

where $\bar{M} = (I - \bar{P} T_c^{-1} \bar{P}^T T_f) S_T^\nu$

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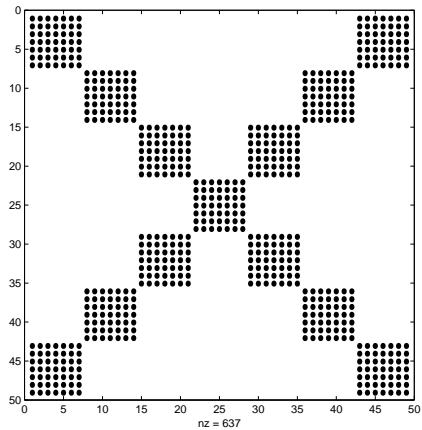
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Q_f is orthogonal:

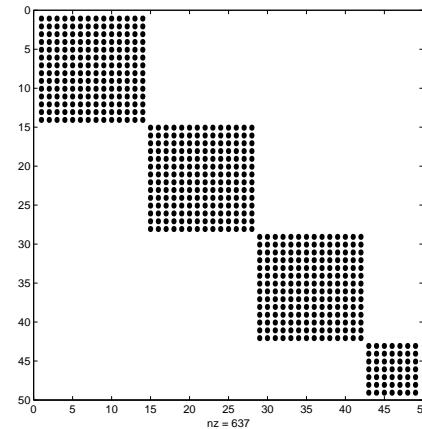
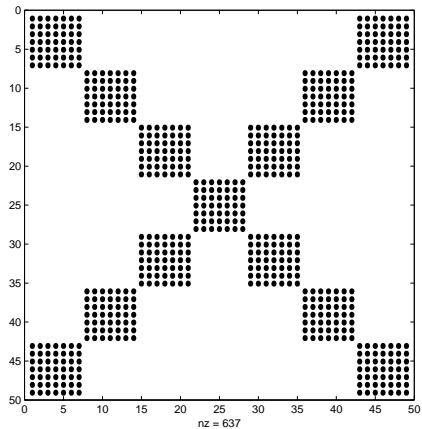
$$\|M\|_2 = \|\bar{M}\|_2$$

Transformed Iteration Matrix (2)



$$\begin{bmatrix} \mathbf{B}_1 & & & & & & \mathbf{C}_1 \\ & \mathbf{B}_2 & & & & & \mathbf{C}_2 \\ & & \mathbf{B}_3 & \mathbf{C}_3 & & & \\ & & & \mathbf{B}_4 & & & \\ & & & & \mathbf{C}_5 & \mathbf{B}_5 & \\ & & & & & & \mathbf{B}_6 \\ \mathbf{C}_6 & & & & & & \mathbf{B}_7 \\ & \mathbf{C}_7 & & & & & \end{bmatrix}$$

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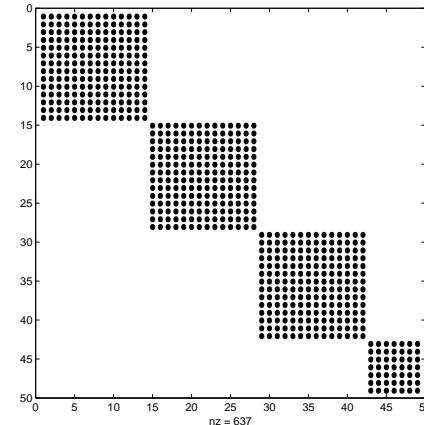
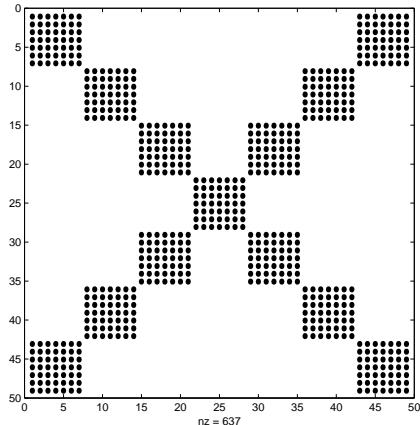


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$$\begin{bmatrix} \mathbf{B}_1 \mathbf{C}_1 & & & & & & \\ & \mathbf{C}_7 \mathbf{B}_7 & & & & & \\ & & \mathbf{B}_2 \mathbf{C}_2 & & & & \\ & & & \mathbf{C}_6 \mathbf{B}_6 & & & \\ & & & & \mathbf{B}_3 \mathbf{C}_3 & & \\ & & & & & \mathbf{C}_5 \mathbf{B}_5 & \\ & & & & & & \mathbf{B}_4 \end{bmatrix}$$

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$$\|\bar{M}\|_2 = \max \left\{ \max_{j=1, \dots, n_c} \left\| \begin{bmatrix} B_j & C_j \\ C_k & B_k \end{bmatrix} \right\|_2, \|B_{N_c}\|_2 \right\}, \quad k = N_f - j$$

The Story So Far...

- $n_f^2 \times n_f^2$ two-grid iteration matrix M
- Fourier transformation converts 2D problem to a set of n_f problems with 1D structure
- $\|M\|_2$ can be found from norms of N_c smaller problems
 n_c of size $2n_f \times 2n_f$, 1 of size $n_f \times n_f$

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- **IDEA:** analyse semiperiodic version of the problem
 n_c of size $2N_f \times 2N_f$, 1 of size $N_f \times N_f$
- gain insight into Dirichlet problem behaviour?

Semiperiodic problem

- B_j, C_j are replaced by periodic versions, e.g.

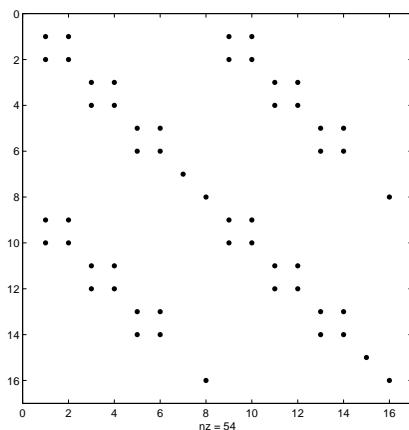
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- B_j^{per}, C_j^{per} become block diagonal with 2×2 blocks

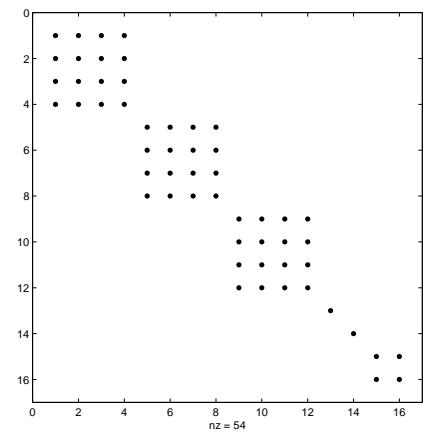
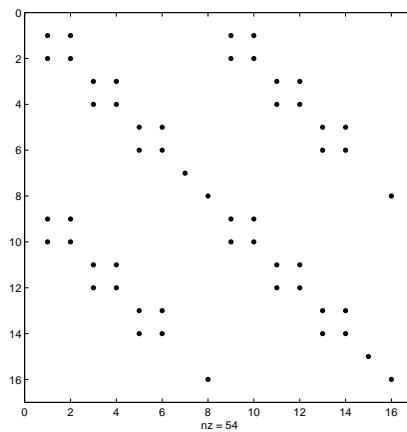


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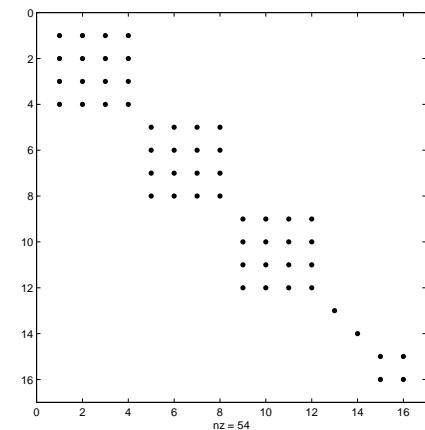
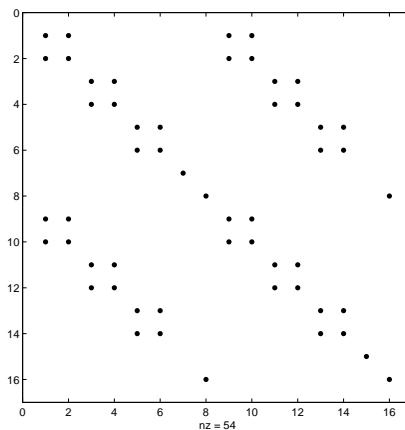


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- 2-norm given by maximum 2-norm of the 4×4 blocks

Analytic result

- with semiperiodic approximation, when $P_h > 1$

$$\|M^{per}\|_2 = \frac{\sqrt{3 + \cos(2\pi h)}}{\sqrt{2}(5^\nu)}$$

independent of ϵ

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$$\|M^{per}\|_2 = \frac{\sqrt{3 + \cos(2\pi h)}}{\sqrt{2}(5^\nu)}$$

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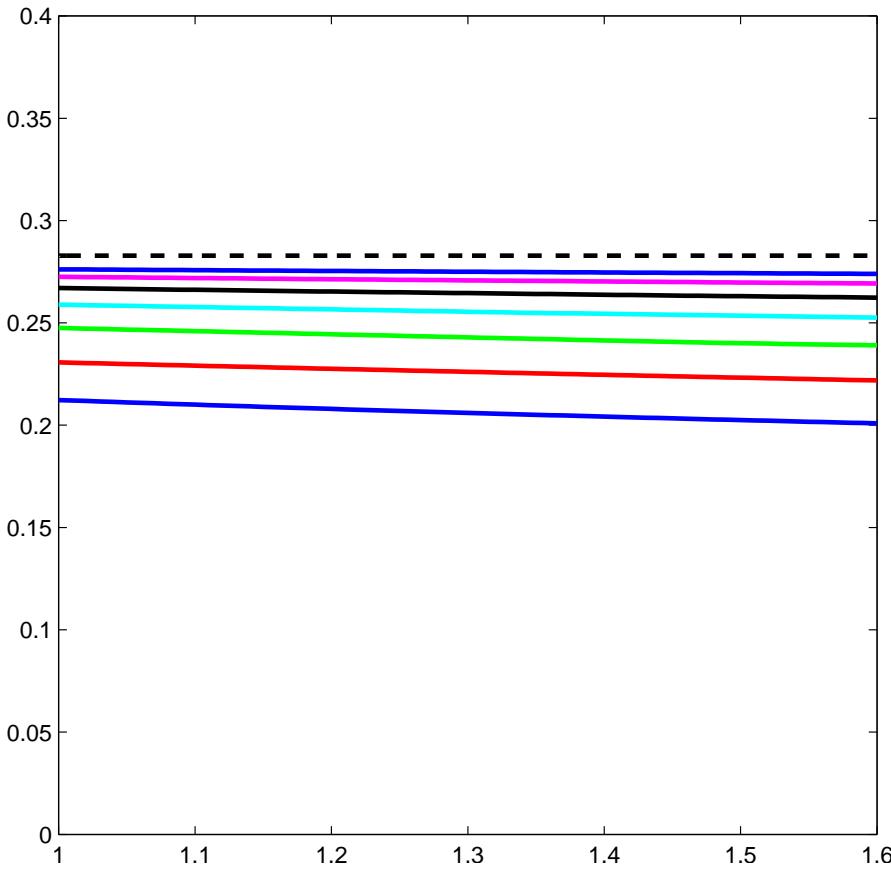
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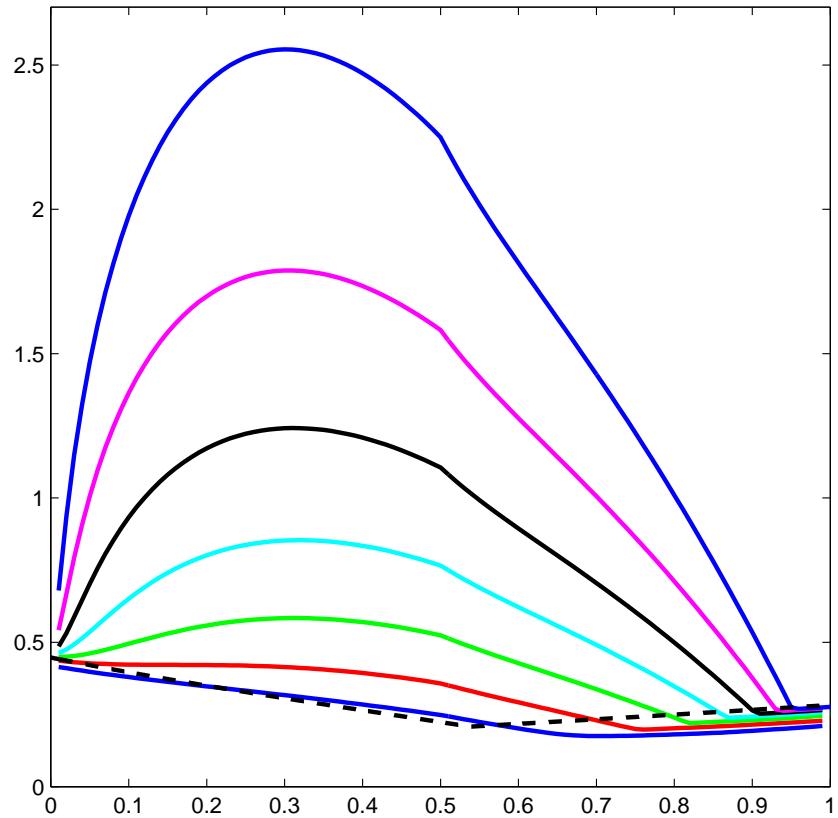
Question: Does this **semiperiodic** analysis correctly predict
Dirichlet problem behaviour?

Model Problem Results (1)



- $\|M\|_2$ vs P_h
- $P_h \geq 1$ only
- semiperiodic: dashed line
- Dirichlet: solid lines
- h fixed for each line
- $h = \frac{1}{8}$ to $h = \frac{1}{512}$
- $\nu = 1$
- semiperiodic: $\frac{\sqrt{2}}{5} \simeq 0.28$
- Dirichlet $\rightarrow \frac{\sqrt{2}}{5}$

Model Problem Results (2)



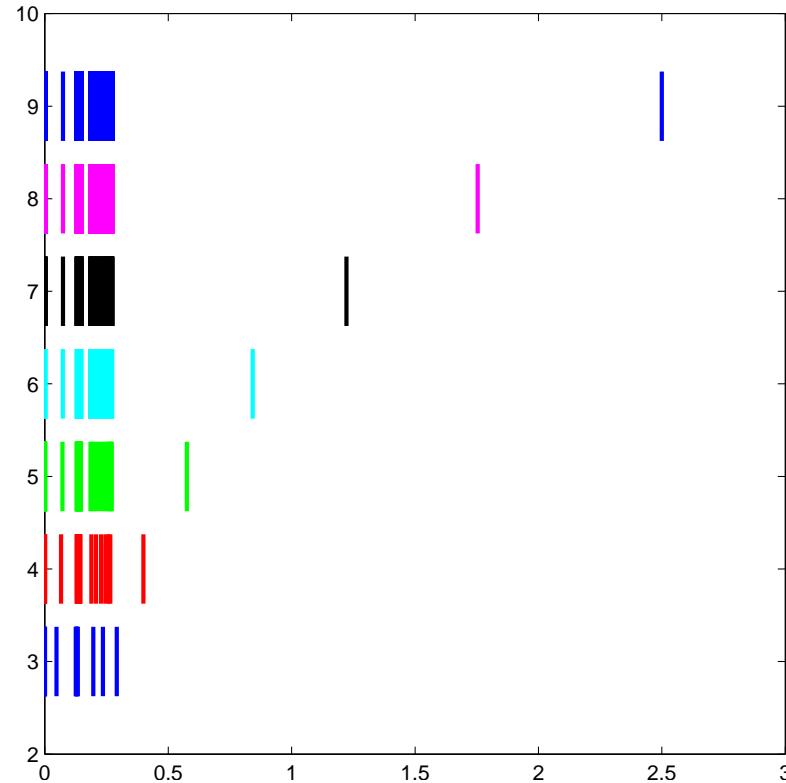
- $\|M\|_2$ vs P_h
- $P_h < 1$ only
- semiperiodic: dashed line
- Dirichlet: solid lines
- h fixed for each line
- $h = \frac{1}{8}$ to $h = \frac{1}{512}$
- $\nu = 1$
- not a good match
- MG may diverge!

Observations

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- for $P_h < 1$, matrix blocks have one ‘bad’ eigenvalue



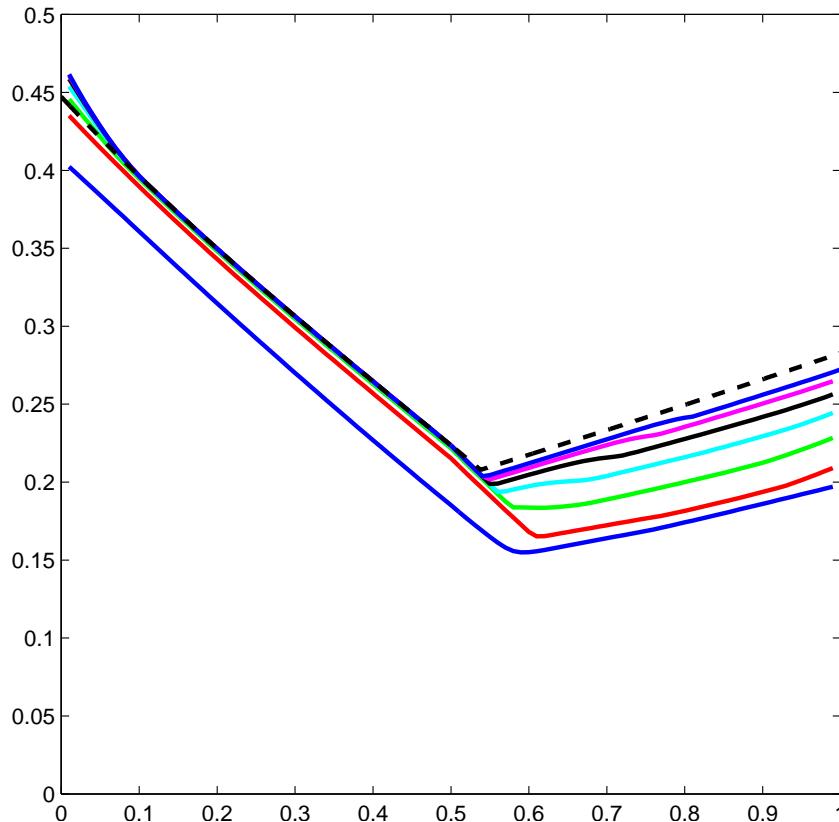
$\sqrt{|\lambda_1(\mathcal{M}_1^*\mathcal{M}_1)|}$ for fixed $P_h = 0.38$

Alternative Bound?

- artificially ‘remove’ this eigenvalue: use $\sqrt{|\lambda_2(M^*M)|}$

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- Dirichlet: $\sqrt{|\lambda_2(M^*M)|}$

MG Iteration Counts

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	ϵ										
h	$\frac{1}{2}$	$\frac{1}{4}$	$\frac{1}{8}$	$\frac{1}{16}$	$\frac{1}{32}$	$\frac{1}{64}$	$\frac{1}{128}$	$\frac{1}{256}$	$\frac{1}{512}$	$\frac{1}{1024}$	$\frac{1}{2048}$
$\frac{1}{4}$	5	5	5	5	5	4	4	3	2	2	2
$\frac{1}{8}$	7	7	6	6	5	5	4	4	3	2	2
$\frac{1}{16}$	7	7	7	6	5	5	5	4	4	3	2
$\frac{1}{32}$	7	7	7	7	6	5	5	4	4	3	3
$\frac{1}{64}$	7	7	7	7	6	5	5	4	4	4	3
$\frac{1}{128}$	7	6	6	6	6	6	5	4	4	4	3

$$P_h < 1$$

$$P_h \geq 1$$

MG-like convergence for any value of P_h

Final Remarks

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 - semiperiodic analysis for approximation matrix norm is representative of Dirichlet problem behaviour for $P_h \geq 1$: for $P_h < 1$, one ‘bad’ eigenvalue again causes trouble.

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 - semiperiodic analysis for approximation matrix norm is representative of Dirichlet problem behaviour for $P_h \geq 1$: for $P_h < 1$, one ‘bad’ eigenvalue again causes trouble.
- Replacing the Dirichlet condition by a Neumann condition on the outflow boundary leads to similar computational results.