

A Moving Mesh Finite Element Method for Modelling Defects in Liquid Crystals

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Motivation

- Liquid crystals may have different **equilibrium** configurations: can force switching between **stable** states by altering applied voltage, magnetic field, boundary conditions, . . .
- Understanding the formation and dynamics of **defects** is important in the design and control of liquid crystal devices.
- Defects typically induce distortion over **very small length scales** as compared to the size of the cell.
- This poses **significant challenges** for standard numerical modelling techniques.
- In this talk we present a finite-element based **adaptive moving mesh** model designed to track defect movement.

Q-tensor representation

- Represent **average** molecular orientation by **symmetric traceless** order tensor

$$\mathbf{Q} = \sqrt{\frac{3}{2}} \left\langle \mathbf{u} \otimes \mathbf{u} - \frac{1}{3} \mathbf{I} \right\rangle$$

with **five** degrees of freedom.

- Represent **Q** using a (non-unique) basis of five linearly-independent tensors, e.g.

$$\mathbf{Q} = \begin{bmatrix} q_1 & q_2 & q_3 \\ q_2 & q_4 & q_5 \\ q_3 & q_5 & -q_1 - q_4 \end{bmatrix}.$$

- **Five** unknowns for PDE model: q_1, q_2, q_3, q_4, q_5 .

- Minimise the free energy

$$F = \int_V F_{bulk}(\mathbf{Q}, \nabla \mathbf{Q}) dv$$

$$F_{bulk} = F_{elastic} + F_{thermotropic} + F_{electrostatic}$$

$$F_{elastic} = \frac{1}{2}L_1(\operatorname{div} \mathbf{Q})^2 + \frac{1}{2}L_2|\nabla \times \mathbf{Q}|^2$$

$$F_{thermotropic} = \frac{1}{2}A(T - T^*) \operatorname{tr} \mathbf{Q}^2 - \frac{\sqrt{6}}{3}B \operatorname{tr} \mathbf{Q}^3 + \frac{1}{4}C(\operatorname{tr} \mathbf{Q}^2)^2$$

$$F_{electrostatic} = -\frac{1}{2}\epsilon_0 \mathbf{E} \cdot \epsilon \mathbf{E} - (\bar{\epsilon} \operatorname{div} \mathbf{Q}) \cdot \mathbf{E}$$

- Solutions with **least** energy are physically relevant: solve **Euler-Lagrange** equations.

Derivation of time-dependent PDEs

- Use a **dissipation function** with viscosity coefficient ν :

$$\mathcal{D} = \frac{\nu}{2} \text{tr} \left[\left(\frac{\partial \mathbf{Q}}{\partial t} \right)^2 \right] = \nu (\dot{q}_1 \dot{q}_4 + \dot{q}_1^2 + \dot{q}_2^2 + \dot{q}_3^2 + \dot{q}_4^2 + \dot{q}_5^2).$$

- Obtain **Q**-tensor PDEs (for $i = 1, \dots, 5$ and $j = 1, 2, 3$):

$$\frac{\partial \mathcal{D}}{\partial \dot{q}_i} = \nabla \cdot \hat{\Gamma}_i - \hat{f}_i,$$

$$(\hat{\Gamma}_i)_j = \frac{\partial F_{bulk}}{\partial q_{i,j}}, \quad q_{i,j} = \frac{\partial q_i}{\partial x_j}, \quad \hat{f}_i = \frac{\partial F_{bulk}}{\partial q_i}.$$

- Combining equations and manipulating terms we can write

$$\frac{\partial q_i}{\partial t} = \nabla \cdot \Gamma_i - f_i, \quad i = 1, \dots, 5.$$

Coupling with electric field

- Additional unknown U such that $\mathbf{E} = -\nabla U$.
- Assuming no free charges, solve the **Maxwell equation** $\nabla \cdot \mathbf{D} = 0$ for electric displacement \mathbf{D} .

SUMMARY

- Final time-dependent physical PDEs (**PPDEs**) are

$$\frac{\partial q_i}{\partial t} = \nabla \cdot \mathbf{\Gamma}_i - f_i, \quad i = 1, \dots, 5,$$

$$\nabla \cdot \mathbf{D} = 0.$$

- 6 PDEs in 6 unknowns ($q_1, q_2, q_3, q_4, q_5, U$)

Adapt PPDEs for mesh movement

- Focus here on **Moving Mesh PDE** model.
Huang and Russell, *Adaptive Moving Mesh Methods*, Springer (2011)
- Define **physical** domain Ω and **computational** domain Ω_c .
- Map $\xi = (\xi, \eta) \in \Omega_c$ to $\mathbf{x} = (x, y) \in \Omega$ using bijective mappings $\mathcal{A}_t : \Omega_c \rightarrow \Omega$ such that
$$\mathbf{x}(\xi, t) = \mathcal{A}_t(\xi).$$
- Define a **mesh velocity**

$$\dot{\mathbf{x}}(\mathbf{x}, t) = \left. \frac{\partial \mathbf{x}}{\partial t} \right|_{\xi} (\mathcal{A}_t^{-1}(\mathbf{x}))$$

and apply the Chain Rule to get

$$\left. \frac{\partial q}{\partial t} \right|_{\xi} = \left. \frac{\partial q}{\partial t} \right|_{\mathbf{x}} + \dot{\mathbf{x}} \cdot \nabla q.$$

Finite elements for the physical PDEs

- PPDEs in computational domain ($i = 1, \dots, 5$):

$$\frac{\partial q_i}{\partial t} \Big|_{\xi} - \dot{\mathbf{x}} \cdot \nabla q = \nabla \cdot \Gamma_i - f_i, \quad \nabla \cdot \mathbf{D} = 0.$$

- Find $q_{ih}(t)$, U_h such that, for test functions v_h ,

$$\frac{d}{dt} \int_{\Omega} q_{ih} v_h \, d\mathbf{x} - \int_{\Omega} (\nabla \cdot (\dot{\mathbf{x}} q_{ih})) v_h \, d\mathbf{x} = \int_{\Omega} \Gamma_{ih} \cdot \nabla v_h \, d\mathbf{x} - \int_{\Omega} f_{ih} v_h \, d\mathbf{x},$$
$$\int_{\Omega} \mathbf{D}_h \cdot \nabla v_h \, d\mathbf{x} = 0.$$

- Non-linear differential algebraic system ($i = 1, \dots, 5$)

$$\frac{d}{dt}(M(t)\mathbf{q}_i(t)) = \mathbf{G}_i(t, \mathbf{q}_i(t), \mathbf{u}(t)), \quad \mathbf{C}(\mathbf{q}_i(t), \mathbf{u}(t)) = \mathbf{0}.$$

Moving Mesh PDEs

- Avoid **mesh crossings** by evolving the inverse mapping

$$\mathcal{A}_t^{-1}(\mathbf{x}) = \boldsymbol{\xi}(\mathbf{x}, t).$$

- Choose mapping **$\boldsymbol{\xi}(\mathbf{x})$** for a fixed t to minimise

$$I[\boldsymbol{\xi}] = \frac{1}{2} \int_{\Omega_t} [(\nabla \boldsymbol{\xi})^T G^{-1}(\nabla \boldsymbol{\xi}) + (\nabla \boldsymbol{\eta})^T G^{-1}(\nabla \boldsymbol{\eta})] \, d\mathbf{x}$$

with 2×2 symmetric positive definite **monitor matrix** G .

- For robustness, evolve mesh via **gradient flow** equations

$$\frac{\partial \boldsymbol{\xi}}{\partial t} = \frac{P}{\tau} \nabla \cdot (G^{-1} \nabla \boldsymbol{\xi}), \quad \frac{\partial \boldsymbol{\eta}}{\partial t} = \frac{P}{\tau} \nabla \cdot (G^{-1} \nabla \boldsymbol{\eta}).$$

- User-specified parameters:

- positive **temporal smoothing** parameter τ ;
- positive **spatial balancing** function $P(\mathbf{x}, t)$.

- Use **Winslow** monitor matrix with **monitor function** $w(\mathbf{x}, t)$:

$$G = \begin{bmatrix} w & 0 \\ 0 & w \end{bmatrix}.$$

- In practice, interchange variable roles in MMPDE to obtain

$$\tau \frac{\partial \mathbf{x}}{\partial t} = P(ax_{\xi\xi} + bx_{\xi\eta} + cx_{\eta\eta} + dx_{\xi} + ex_{\eta}).$$

$$a = \frac{1}{w} \frac{x_{\eta}^2 + y_{\eta}^2}{J^2}, \quad b = -\frac{2}{w} \frac{(x_{\xi}x_{\eta} + y_{\xi}y_{\eta})}{J^2}, \quad c = \frac{1}{w} \frac{x_{\xi}^2 + y_{\xi}^2}{J^2},$$

$$d = \frac{1}{(wJ)^2} [w_{\xi}(x_{\eta}^2 + y_{\eta}^2) - w_{\eta}(x_{\xi}x_{\eta} + y_{\xi}y_{\eta})],$$

$$e = \frac{1}{(wJ)^2} [-w_{\xi}(x_{\xi}x_{\eta} + y_{\xi}y_{\eta}) + w_{\eta}(x_{\xi}^2 + y_{\xi}^2)].$$

Overview of full algorithm

Set an initial uniform mesh Δ_N^0 . Set the initial guess \mathbf{q}_i^0 .

Select an initial Δt^0 . Set $n = 0$.

while ($t^n < t^{\max}$);

 Evaluate monitor function at time t^n .

 Integrate **MMPDE** forward in time to obtain new grid Δ_N^{n+1} .

 Integrate **PPDEs** forward using SDIRK2 to obtain $\mathbf{q}_i^{n+1}, \mathbf{u}^{n+1}$.

$n := n + 1$.

end while.

Choice of monitor function

- **BM2**: Based on **second-order partial derivatives** of $\mathcal{T}(\mathbf{x}, t)$:

$$w(\mathcal{T}(\mathbf{x}, t)) = \alpha(\mathbf{x}, t) + \left(\sqrt{\left(\frac{\partial^2 \mathcal{T}}{\partial x^2}\right)^2 + 2\left(\frac{\partial^2 \mathcal{T}}{\partial x \partial y}\right)^2 + \left(\frac{\partial^2 \mathcal{T}}{\partial y^2}\right)^2} \right)^{\frac{1}{m}}$$

- Scaling parameters α and m regulate **mesh clustering**.
- Input function based on a direct invariant measure of **biaxiality**

$$\mathcal{T}(\mathbf{x}, t) = \left[1 - \frac{6 \operatorname{tr}(\mathbf{Q}^3)^2}{\operatorname{tr}(\mathbf{Q}^2)^3} \right]^{\frac{1}{2}}.$$

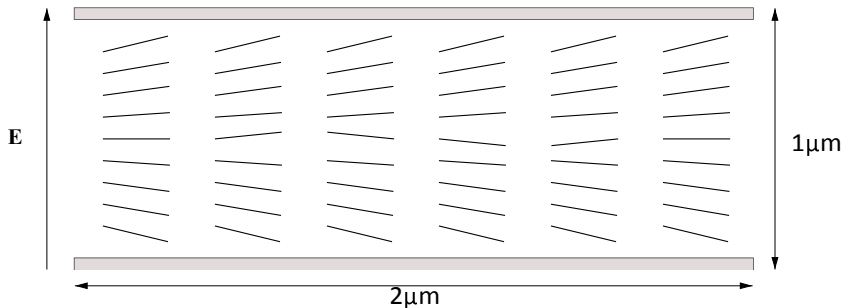
- This has an extremum at the centre of a defect and **varies rapidly** in the immediate neighbourhood of the defect centre.

Test problem: 2D Pi-cell

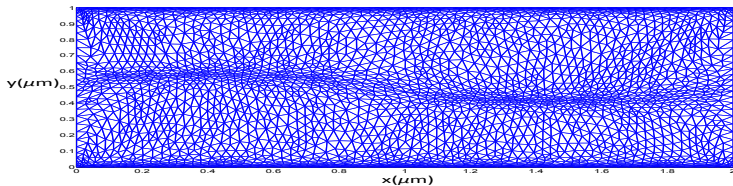
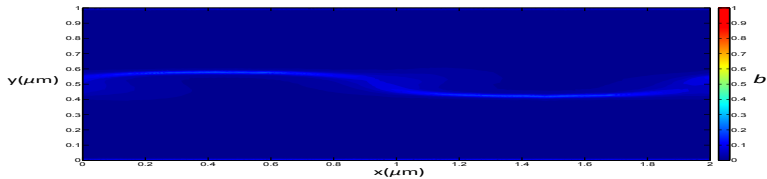
- Two-dimensional **Pi-cell** geometry.
Zhang, Chung, Wang and Bos, *Liquid Crystals* 34(2), 2007
- Electric field applied parallel to the cell thickness.
- Inhomogeneous transition mediated by the nucleation of **defect pairs** moving and annihilating each other.
- Initial director angle across cell centre follows $\sin(2\pi x/p)$ for cell width p at $t = 0$ for one time step.
- Use **triangular grid** with **quadratic** basis functions for PPDEs, **linear** basis functions for MMPDE.

Pi-cell geometry

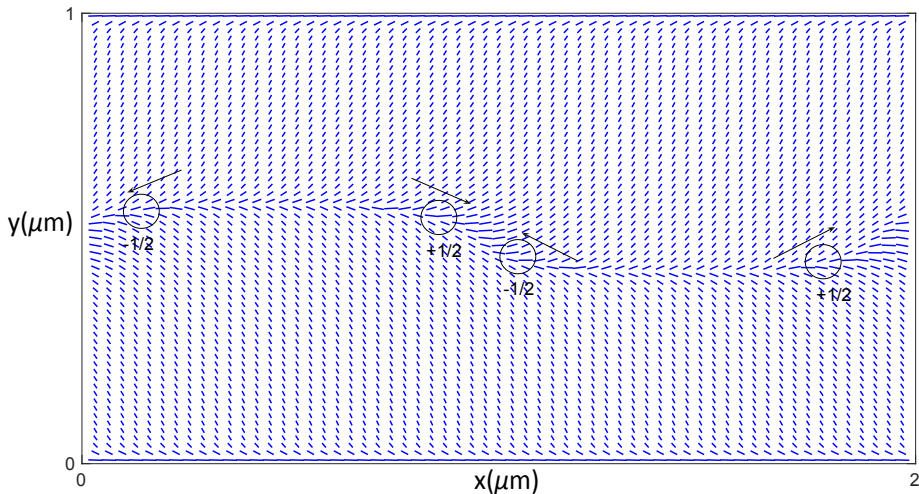
- Pre-tilt angle $\theta = \pm 6^\circ$ at boundaries.
- Electric field strength $18V\mu\text{m}^{-1}$.



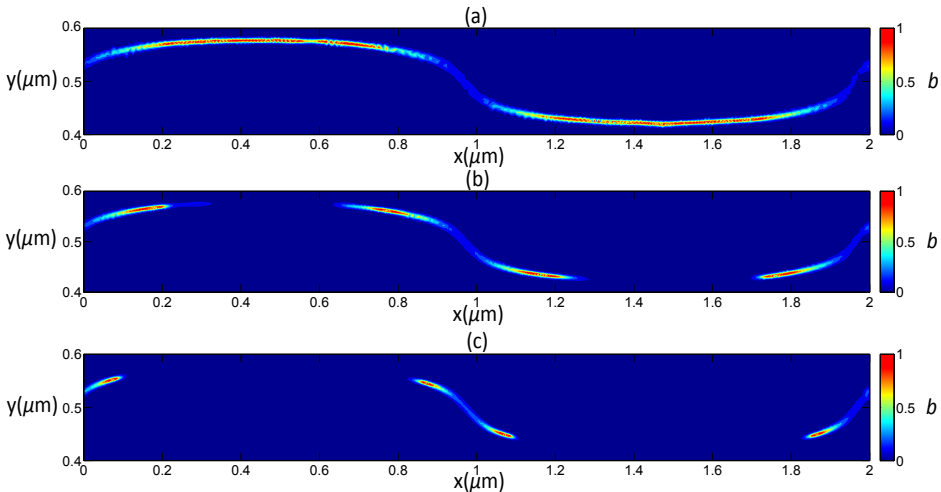
Biaxiality and mesh after $12\mu\text{s}$



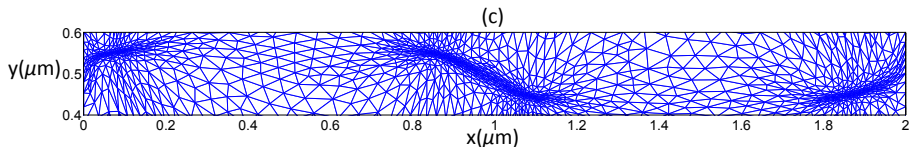
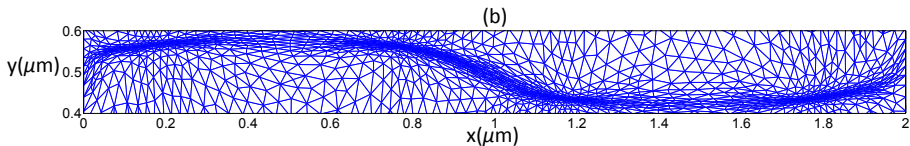
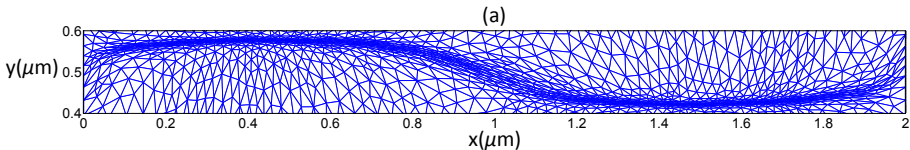
Director field after $15.5\mu\text{s}$



Biaxiality after (a) $15.5\mu\text{s}$ (b) $16\mu\text{s}$ and (c) $17\mu\text{s}$



Adaptive mesh after (a) $15.5\mu\text{s}$ (b) $16\mu\text{s}$ and (c) $17\mu\text{s}$



Summary and future work

- New efficient **moving mesh method** for **Q**-tensor models of liquid crystal cells.
- Found **biaxiality** to be a good choice for the monitor input function.
- Demonstrated **optimal** spatial convergence for a model of a static $+1/2$ defect.
- Method resolved the movement and core details of **defects** (including **creation** and **annihilation**) in a time-dependent Pi-cell problem.
MacDonald, Mackenzie and Ramage, JCP:X 8, 2020
- Future challenges involve the extension to more irregular geometries (e.g. the ZBD) and three dimensions.